Playing with a photon

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Matter is made of atoms (image : Wikipedia)

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Matter is made of atoms (image : Wikipedia)



Light is made of photons (image : CNRS)

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How to keep light for a while





Prepare and probe a delocalized state of one photon stored in two cavities :

$$rac{1}{\sqrt{2}}\left(\left| 10
ight
angle + \left| 01
ight
angle
ight).$$

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Designing highly reflective mirrors

- Microwave frequency : ν(TEM₉₀₀) = 51.099 GHz
- Long lifetime $T_{cav} =$ some tens of milliseconds !
- Small mode volume $V \simeq 700 \text{ mm}^3$

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$$\mathcal{E}_0 = \sqrt{\frac{\hbar\omega_c}{2\varepsilon_0 V}} \simeq 1.5 \text{ mV/m}$$







Boxes to trap light









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Circular Rydberg atoms

$$|e\rangle \equiv |51c\rangle \frac{1}{\nu} = 51.099 \text{ GHz}$$

- ⁸⁵Rb
- Rydberg atoms : large principal quantum number $n \simeq 50$
- Circular Rydberg : maximal angular momentum : l = |m| = n 1
- Long lifetime $T_{\rm at} \simeq 30 \ {\rm ms}$
- Microwave transitions
- Large dipole : $d=1776|q|a_0$ for |50c
 angle
 ightarrow|51c
 angle

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Preparing atomic samples



- Paquets with typically 0.1 atoms
- Velocity 250 m/s
- Interaction with cavity controled by V
- Ionization detector D : detects $|e\rangle$ or $|g\rangle$

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State superposition

Goal : we want to understand a notation such as $|\pm_x\rangle = (|e\rangle \pm |g\rangle)/\sqrt{2}$. Note that :

$$P_e = \langle e|+_x \rangle = rac{1}{2},$$

 $P_g = \langle g|+_x \rangle = rac{1}{2}.$

But if we change the measurement basis :

$$P_{+} = \langle +_{x} | +_{x} \rangle = 1,$$

$$P_{-} = \langle -_{x} | +_{x} \rangle = 0.$$

What does it mean to change the measurement basis?

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The Bloch sphere



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The Bloch sphere



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The Bloch sphere



Most general state (up to a global phase) :

$$|\psi\rangle = \cos\theta \,|e\rangle + \sin\theta e^{i\varphi} \,|g\rangle = \begin{pmatrix} \cos\theta \\ \sin\theta e^{i\varphi} \end{pmatrix}.$$



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Atom-cavity interaction : The Jaynes-Cummings Hamiltonian

$$\hat{H}_{\rm JC} = \frac{\hbar\omega_{\rm at}}{2}\hat{\sigma}_z + \hbar\omega_c \left(\hat{N} + \frac{1}{2}\right) + \frac{\hbar\Omega_0}{2} \left(\hat{a}\hat{\sigma}_+ + \hat{a}^{\dagger}\hat{\sigma}_-\right)$$
$$\Omega_0 = \frac{2d\mathcal{E}_0\vec{\epsilon}_a^*\cdot\vec{\epsilon}_c}{\hbar} \simeq 2\pi\cdot 50 \text{ kHz} >> T_{\rm at}^{-1}, T_{\rm cav}^{-1}$$

Interaction of one atom with one photon : Rabi oscillations

- *Ĥ*_{jc} commutes with the number of excitations and can be separated in terms *Ĥ_n* with *n* excitations.
- For an atom initially in $|e\rangle$ with no photon, one can restrict to the subspace of states $\{|e, 0\rangle, |g, 1\rangle\}.$





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Preparation of an entangled state



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Reading of the entangled state



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The role of the quantum phase



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The role of the quantum phase



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To absorb or not to absorb



Interferometric signal



- $\bullet\,$ Signature of the non-local coherence between $|10\rangle$ and $|01\rangle$
- Signal reveals the frequency difference between the two cavities

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Reconstructed density matrix



- Preparation and probe of a state |10
 angle+|01
 angle, with tomography;
- The non-local coherences play a major role in the sensitivity of the state;
- The photon is sensitive to the frequency beat of the cavities : it's in both of them at the same time !

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 angle, with tomography;
- The non-local coherences play a major role in the sensitivity of the state;
- The photon is sensitive to the frequency beat of the cavities : it's in both of them at the same time !
- Perpective : Preparation of $|20\rangle+|02\rangle$

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Thank you for your attention !

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The density matrix in a nutshell

$$\begin{split} |e\rangle &\to |e\rangle \langle e| = \begin{pmatrix} 1\\ 0 \end{pmatrix} \begin{pmatrix} 1 & 0 \end{pmatrix} = \begin{pmatrix} e & g\\ 1 & 0\\ 0 & 0 \end{pmatrix} \\ |+\rangle &\to |+\rangle \langle +| = \frac{1}{2} \begin{pmatrix} 1\\ 1 \end{pmatrix} \begin{pmatrix} 1 & 1 \end{pmatrix} = \frac{1}{2} \begin{pmatrix} 1 & 1\\ 1 & 1 \end{pmatrix} \\ |\psi_{\varphi}\rangle &\equiv (|e\rangle + e^{i\varphi} |g\rangle)/\sqrt{2} \to \frac{1}{2} \begin{pmatrix} 1 & e^{-i\varphi}\\ e^{i\varphi} & 1 \end{pmatrix} \end{split}$$

Moere generally, every matrix that fullfills :

$$egin{aligned} & \end{aligned} \Gamma r(\hat{
ho}) &= 1, \ & \hat{
ho}^\dagger &= \hat{
ho}, \ & \hat{
ho} &> 0, \end{aligned}$$

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is a density matrix.

Diagonalizing some density matrices

Diagonalization of
$$\hat{\rho}_{\varphi} \equiv |\psi_{\varphi}\rangle \langle \psi_{\varphi}| = \frac{1}{2} \begin{pmatrix} 1 & e^{-i\varphi} \\ e^{i\varphi} & 1 \end{pmatrix}$$
:
 $\hat{\rho}_{\varphi} |\psi_{\varphi}\rangle = |\psi_{\varphi}\rangle \langle \psi_{\varphi}|\psi_{\varphi}\rangle = |\psi_{\varphi}\rangle,$
 $\hat{\rho}_{\varphi} |\psi_{-\varphi}\rangle = |\psi_{\varphi}\rangle \langle \psi_{\varphi}|\psi_{-\varphi}\rangle = 0.$

Hence,

$$\begin{split} \psi_{\varphi} & \psi_{-\varphi} \\ \hat{\rho}_{\varphi} &= \hat{P} \ \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix} \hat{P}^{\dagger}. \end{split}$$
Diagonalization of $\tilde{\hat{\rho}}_{\varphi} \equiv \frac{1}{2} \begin{pmatrix} 1 & ce^{-i\varphi} \\ ce^{i\varphi} & 1 \end{pmatrix} = \frac{1+c}{2} \hat{\rho}_{\varphi} + \frac{1-c}{2} \hat{\rho}_{-\varphi}, \text{ where } (c \in [-1,1]): \end{split}$

$$ilde{
ho}_{arphi} = \hat{P} egin{pmatrix} rac{1+c}{2} & 0 \ 0 & rac{1-c}{2} \end{pmatrix} \hat{P}^{\dagger}.$$

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• Born's rule :

$$p(m) = \left| \langle m | \psi
angle
ight|^2 o p(m) = \langle m | \, \hat{
ho} \, | m
angle$$

• Average :

$$\langle A
angle = {
m Tr}(\hat{
ho}\hat{A})$$

• Simple expression for $\hat{\rho}$:

$$\hat{\rho} = \frac{1}{2} \left(\mathbbm{1} + \vec{r} \cdot \hat{\vec{\sigma}} \right)$$



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